C–C versus C–H Bond Activation of Alkynes by Early Second-Row Transition Metal Atoms

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Introduction

The oxidative addition of covalent bonds to transition metal centers, M + R–R′ → R–M–R′, plays a significant role in many catalytic and bond-activation processes.1–2 For reactions with hydrocarbons, two oxidative addition processes may compete: activation of a carbon–hydrogen (C–H) bond or activation of a carbon–carbon (C–C) bond.1–10 For reactions of transition metal cations with small molecules in the gas phase, the competition between these processes has been studied quite extensively.6,7 However, for neutral transition metal atoms, which are in general much less reactive with hydrocarbon molecules,7 C–C bond activation is relatively uncommon.8–10

A number of complementary approaches have been taken in theoretical studies of transition metal–hydrocarbon reactions. In one approach, detailed theoretical calculations were carried out for reactions of one or a small number of realistic transition metal systems containing ligands.11 An alternative approach involved studying trends in reactivity of a large number of closely related model systems.3,4,12 For example, the reactions of bare neutral transition metal atoms with simple hydrocarbons such as CH4 were studied theoretically for all three rows of the periodic table.12 The activation of C–H bonds by transition metal atoms involves the donation of electron density from the C–H σ-bonding orbital into metal acceptor orbitals of σ-symmetry, with concurrent back-donation from occupied metal d orbitals into the σ*-antibonding orbital. Metal s^1d^1 electronic configurations were found to be optimal for reducing reactant repulsion upon approach to the transition state (TS), while the formation of two covalent bonds in the C–M–H intermediate was most favorable for low-spin s^1d^1 electronic configurations.3,7 For neutral second-row transition metal atoms, the ground electronic states typically have s^2d^2 configurations.7,13 Therefore, barriers for C–H and C–C insertion result from the avoided crossing between diabatic curves correlating to the ground-state atom and the atomic state(s) leading to the reaction.4,7

It has been noted previously that although the C–H bond energies increase going from ethane (432 kJ/mol) to ethene (459 kJ/mol) to ethyne (549 kJ/mol),14 the potential energy barrier heights for C–H insertion by second-row transition metal atoms actually decrease.3 Since insertion requires favorable orbital overlap between the metal and the C–H σ-bond at the transition state, steric considerations likely play an important role in determining barrier heights.3 It also has been known for many years that alkenes and alkynes can form strongly bound (120−250 kJ/mol) π-association complexes between the metal and the C–C double and triple bonds.3,4 Alkanes, on the other hand, can only form weakly bound (<8 kJ/mol) σ-complexes prior to C–H insertion.11 Although the ability of alkenes and alkynes to form complexes has been thought to play an important role in reducing barriers for reaction,3,4,15 in some cases, evidence has been presented that complex formation does not necessarily precede C–H insertion.16

We previously have studied the reactions of early second-row transition metal atoms with various hydrocarbons. For saturated hydrocarbons (methane, ethane, and propane), only C–H bond activation is observed, yielding dehydrogenation products (e.g., MCH2 + H2, MC2H4 + H2, etc.), or in some cases, formation of the metal dihydride, HMH.17–19 These results are consistent with ab initio calculations that find C–C bond-activation barriers to be larger than C–H bond-activation barriers by 40–80 kJ/mol for neutral metal atom reactions with alkanes.3,4 Reactions of ground-state Y, Zr, Nb, and Mo* with propene also yield dehydrogenation products.10,20 For Y + propene, C–C bond activation was also observed, producing YCH2 + C2H4.10 However, the reaction does not involve insertion of Y into the sp^3–sp^2 C–C bond but instead involves π-complex formation followed by β–C–H bond activation, with subsequent H migration to form a four-membered ring intermediate (metallocyclobutane), which then decays to give YCH2...
+ C₂H₂ products. A similar mechanism was proposed for YCH₃ formation from reactions of Y with four butene isomers.

In reactions of Y + C₃H₂, C–H activation leading to the formation of YC≡CH + H₂ was found to be dominant. As in the studies of reactions involving alkenes, strong nonreactive wide-angle scattering of metal atoms was observed, demonstrating that a substantial fraction of initially formed π-complexes decayed back to reactants. At sufficiently high collision energies, the formation of YCH + H also was observed. From the observed threshold behavior for this reaction, we concluded that \( D_0(\text{Y–CCH}) = 461.1 \pm 8.4 \text{ kJ/mol} \).

In this study, we observed a number of competing pathways for reactions of second-row transition metal atoms with propyne and 2-butyne. For ground-state Y atoms, the following reactions were observed:

\[
\begin{align*}
Y + \text{CH}_2\text{CCH} & \rightarrow \text{YCH} + \text{CH}_3 & \Delta E \approx +60 \text{ kJ/mol} \\
& \rightarrow \text{YCH}_2 + \text{H} & \Delta E \approx -130 \text{ kJ/mol} \\
Y + \text{CH}_2\text{CCCH} & \rightarrow \text{YCH}_3 + \text{CH}_3 & \Delta E \approx +70 \text{ kJ/mol} \\
& \rightarrow \text{Y(CH}_3)_2 + \text{H} & \Delta E \approx -100 \text{ kJ/mol} \\
& \rightarrow \text{YCH}_2 + \text{CH}_3 & \Delta E \approx -140 \text{ kJ/mol}
\end{align*}
\]

The goal of the present study was to explore the competition between C–H and C=C activation in reactions of alkynes containing three or four carbon atoms with Y, Nb, Zr, and Mo.

**Experimental Procedures**

The reactions with propyne and 2-butyne were investigated using a rotatable source crossed molecular beams apparatus. Neutral metal atoms were produced by laser vaporization and entrained in a supersonic expansion of an inert carrier gas (He or Ne). The resulting atomic electronic-state populations were characterized by laser induced fluorescence (LIF) excitation spectroscopy. For Y, Zr, and Mo, only the ground electronic states (a°D, a1F, and a1S, respectively) were observed. For Nb, small populations of excited a4F and a4P states were observed in addition to the ground a1D state.

Photodepletion studies were performed for Nb to investigate the role of the ground a1D electronic state in product formation. Frequency-doubled dye laser radiation (6 mJ) was introduced upstream of the interaction region to deplete the Nb beam by 1 + 1 REMPI of ground-state atoms. To maintain alignment while the molecular source was rotated, the depletion laser was aligned along the axis of rotation and periscoped 1 cm upstream of the reaction zone using a 0.75 in. thick window tilted at 45° and mounted to the rotatable source. Because the detector lenses were placed at +100 V, metal ions or the ionic products of their reactions were unable to enter the detector. The z2S1/2 \( \rightarrow \) a1D3/2 and y2P3/2 \( \rightarrow \) a1D5/2 transitions were simultaneously pumped near 358 nm. The similarity between the fractional depletion of the Nb ground-state reactant beam to the depletion of the product signal indicated that the products resulted primarily from reactions of ground-state Nb a1D atoms.

**Results**

**M + Propyne. Nonreactive Scattering.** Nonreactive scattering of M from propyne was recorded for lab angles \( \theta = 5–50^\circ \), where \( \theta = 0^\circ \) corresponds to the M atomic beam. Figure 1 shows representative TOF spectra recorded using REMPI for the ground states of Y through Mo at a collision energy \( (E_{\text{coll}}) \) of 99 kJ/mol. For Mo (a1S), the solid lines in Figure 1 were generated using a forward peaking CM angular distribution, \( T(\theta) \) (Figure 2a). The Mo (a1S) data were simulated using two weighted \( T(\theta) \) distributions (dotted and dashed lines in Figure 2a) associated with two separate \( P(E) \) distributions. For Y, Zr, and Nb, a slow peak \( (~190 \mu s) \) became more significant at wider angles, and it was necessary to use a symmetric, forward-backward peaking \( T(\theta) \) as shown in Figure 2b in addition to a forward peaking \( T(\theta) \) similar in shape to that in Figure 2a. Each \( T(\theta) \) distribution was associated with a separate \( P(E) \) distribution. The solid-line fits shown for Y, Zr, and Nb in Figure 1 are the sum of the contributions from the forward peaking \( T(\theta) \) (dotted lines) and from the forward-backward peaking \( T(\theta) \) (dashed-dotted lines). Figure 3 shows the TOF spectra for Mo nonreactive scattering, obtained with the pump laser turned on. As shown in the lower part of Figure 3, the forward scattered component decreases, and the symmetric component increases when the laser is turned on. The contribution from the symmetric component decreases in the order \( Y(26\%) > Mo(20\%) > Nb-(17\%) > Zr(8\%) > Mo(0\%) \).

**H₂ Elimination.** The reaction M + propyne \( \rightarrow \text{MC}_3\text{H}_2 + \text{H}_2 \) was observed for all ground electronic-state atoms studied and for Mo (a1S). Figure 4 shows laboratory angular distributions for MC₃H₂ products \( (E_{\text{coll}} \approx 50 \text{ kJ/mol}) \) for each metal. Photodepletion of Nb (a1D) from the atomic beam was correlated to depletion of the NbC₃H₂ product signal at several lab angles, unambiguously identifying the reaction as occurring primarily from ground-state Nb atoms.

The H₂ elimination data were simulated using the \( P(E) \) distributions included in Figure 5 (top panel) with a nearly isotropic \( T(\theta) \) distribution that was symmetric with respect to \( \theta = 90^\circ \). The translational energy released to products was largest for the reaction with Y and identical for reactions with Zr, Nb, and Mo. Comparing Mo (a1S) to Mo (a1S) indicated that a slightly larger amount of translational energy was deposited in products from Mo* reactions. As shown in Figure 5 (bottom panel), the \( P(E) \) values were divided by the total...
amount of available energy. We observed a greater translational energy release for Y as compared to Zr and Nb (e.g., \( \langle f(T) \rangle \approx 0.29 \) as compared to \( \langle f(T) \rangle \approx 0.25 \)).

CH$_3$ Elimination. The reaction M + propyne \( \rightarrow \) MC$_2$H + CH$_3$ was only observed for Y. The TOF spectra for YC$_2$H products obtained at \( E_{\text{coll}} = 99 \) kJ/mol are shown in Figure 6 and have been corrected for a minor contribution from fragmentation of YC$_3$H$_2$ (0.7% of the m/e 127 signal). The solid-line fits were generated using the CM angular distribution included in Figure 7. The average CM translational energy release was 22 kJ/mol, and the T(\( \Theta \)) distributions peaked in the forward and backward directions with \( T(0^\circ)/T(90^\circ) = 5.5 \). Assuming equal 157 nm photoionization cross-sections for YC$_2$H and YC$_3$H$_2$,\(^{9,10} \) the product branching ratio, \( \phi_{\text{YC}_2\text{H}}/\phi_{\text{YC}_3\text{H}_2} \), was 0.11:1.00. The lab angular distributions for these two competing product channels for Y + propyne are shown in Figure 8a.

Competition between CH$_3$ and H$_2$ elimination was monitored as a function of \( E_{\text{coll}} \) (Figure 8). At the lowest \( E_{\text{coll}} \) studies (50 and 54 kJ/mol), a weak signal recorded at m/e 114 was identical in form to the m/e 127 signal, at 0.7% of the intensity. We concluded that the m/e 114 signal at this \( E_{\text{coll}} \) value was due to fragmentation of YC$_3$H$_2$. The real YC$_2$H signal was observed at \( E_{\text{coll}} \geq 63 \) kJ/mol. While simulation of all experimental data required convolution over the \( E_{\text{coll}} \) distribution (Figure 9, top panel), YC$_2$H data obtained near this threshold required a step-function cutoff to accurately fit the width and arrival times for the YC$_2$H TOF spectra.\(^{21} \) Thus, only collisions at energies above the threshold were able to form products. At the lowest \( E_{\text{coll}} \) where CH$_3$ elimination was observed, a cutoff value of 63 kJ/mol was used. Further experiments near this threshold collision energy found that as \( E_{\text{coll}} \) increased, the best-fit energy threshold shifted to higher energies. For example, at \( E_{\text{coll}} = 79 \) kJ/mol, the best-fit cutoff was 67 ± 4 kJ/mol. This behavior is similar to that observed in the Y + C$_2$H$_2$ reaction.\(^{21} \) The maximum translational energy released for YC$_2$H at each \( E_{\text{coll}} \)
value was consistent with this threshold (see, i.e., Figure 7). A plot of the product branching ratio as a function of \(E_{\text{coll}}\) is shown in Figure 9 (lower panel).

**Y + Allene.** The reaction of \(Y\) with allene (CH\(_2\)=C\(=\)CH\(_2\)), a structural isomer of propyne with nearly the same enthalpy of formation, was studied at a collision energy \(E_{\text{coll}}\) \(=\) 84 kJ/mol. The formation of YC\(_3\)H\(_2\)+CH\(_3\) as well as YC\(_3\)H\(_2\)+H\(_2\) was observed. The branching ratio and translational energy disposal were found to be very similar to \(Y + \text{propyne}\). The nonreactive scattering of \(Y + \text{CH}_2=\text{C}==\text{CH}_2\) was also studied. As in the case of \(Y + \text{propyne}\), strong wide-angle scattering of \(Y\) was observed, indicating that a large fraction of collisions formed \(\pi\)-association complexes.

**M + 2-Butyne.** \(Y (\alpha^D) + 2\)-Butyne. The reaction of \(Y\) with 2-butyne was investigated at \(E_{\text{coll}}\) \(=\) 69 and 116 kJ/mol. Three products were detected: YC\(_4\)H\(_4\) \((m/e\) 141), YC\(_3\)H\(_3\) \((m/e\) 128), and YC\(_3\)H\(_2\) \((m/e\) 127). Lab angular distributions for competing channels recorded under identical conditions at \(E_{\text{coll}}\) \(=\) 69 kJ/mol are shown in Figure 10; those at the higher \(E_{\text{coll}}\) values looked similar and are not shown. The optimized CM translational energy release distribution, \(P(E)\), and CM angular distribution, \(T(\theta)\), for each product channel are shown in Figure 11. The \(T(\theta)\) distributions were symmetric with respect to \(\theta = 90^\circ\) and peaked in the forward and backward directions. Using the binding energy for YC\(_3\)H\(_2\) derived for the \(Y + \text{propyne}\) reaction, the average fraction of total available energy released as translational energy for YC\(_3\)H\(_2\)+CH\(_4\) was \(\langle f\rangle (\text{YC}_3\text{H}_2) = 0.32\).

The lab angular distributions depicted in Figure 10 represent the same number of laser shots for each product channel. Methane elimination comprised 19\% of the product signal at \(E_{\text{coll}}\) \(=\) 69 kJ/mol (Table 1). The relative amount of CH\(_3\) elimination increased from 2 to 15\% at the higher collision energy.

**Zr (\(\alpha^F\)) + 2-Butyne.** Two channels were observed: ZrC\(_4\)H\(_4\)+H\(_2\) and ZrC\(_3\)H\(_2\)+CH\(_4\). Laboratory angular distributions recorded at \(m/e\) 142 (\(^{90}\)ZrC\(_4\)H\(_4\)) and \(m/e\) 128 (\(^{90}\)ZrC\(_3\)H\(_2\)) are shown in Figure 12. The product branching ratio, \(\phi_{ZrC_3H_2}/\phi_{ZrC_4H_4}\), was 0.14:1.00 (Table 1). While the lab angular distribution for
H₂ elimination appears qualitatively similar to that for YC₄H₄, the CH₄ elimination angular distribution peaked near \( \Theta_{CM} \), in contrast to the broad distribution observed for YC₃H₂ (Figure 10). In simulating the data for this product channel, the \( T(I) \) distribution was similar to that for Y, but the \( P(E) \) distribution was distinct. Using the values for MC₃H₂ binding energies from the H₂ elimination channels for propyne, the \( P(E) \) values for CH₄ elimination from 2-butyne were converted to \( P(f_T) \) (Figure 13). The average value of energy released as translational energy, \( \langle f_T(ZrC₃H₂) \rangle \), was 0.23.

Figure 8. Lab angular distributions for YC₃H₂ (circles) and YC₂H (squares) products from the Y + propyne reaction recorded at \( E_{coll} = \) (a) 99, (b) 80, (c) 75, (d) 69, and (e) 56 kJ/mol. Product branching ratio, \( \phi_{YC3H2}/\phi_{YC2H} \), included in the upper right corner of each panel. Distributions for YC₂H products were multiplied by a factor of 20 at each collision energy for presentation. Arrows denote center of mass angle for each system.

Figure 9. Top: collision energy distributions a–e associated with Y + propyne data shown in Figure 8, with measured threshold for CH₃ elimination. Bottom: product branching ratio, \( \phi_{YC3H2}/\phi_{YC2H} \), vs \( E_{coll} \).

H₂ elimination appears qualitatively similar to that for YC₄H₄, the CH₄ elimination angular distribution peaked near \( \Theta_{CM} \), in contrast to the broad distribution observed for YC₃H₂ (Figure 10). In simulating the data for this product channel, the \( T(I) \) distribution was similar to that for Y, but the \( P(E) \) distribution was distinct. Using the values for MC₃H₂ binding energies from the H₂ elimination channels for propyne, the \( P(E) \) values for CH₄ elimination from 2-butyne were converted to \( P(f_T) \) (Figure 13). The average value of energy released as translational energy, \( \langle f_T(ZrC₃H₂) \rangle \), was 0.23.

Figure 10. Lab angular distributions for YC₄H₄ (circles), YC₃H₃ (squares), and YC₃H₂ (triangles) products from the Y + 2-butyne reaction recorded at \( E_{coll} = 69 \) kJ/mol. Distributions have been normalized to the same number of laser shots. Solid-line fits generated using CM distributions shown in Figure 11. Arrows denote center of mass angle for each system.

TABLE 1: M + 2-Butyne Product Branching Ratios

<table>
<thead>
<tr>
<th>metal</th>
<th>( E_{coll} ) kJ/mol</th>
<th>( \phi_{YC4H4} )</th>
<th>( \phi_{YC3H3} )</th>
<th>( \phi_{YC3H2} )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Y(( a^2D ))</td>
<td>69</td>
<td>1.00</td>
<td>0.24</td>
<td>0.02</td>
</tr>
<tr>
<td>Y(( a^2D ))</td>
<td>116</td>
<td>1.00</td>
<td>0.16</td>
<td>0.20</td>
</tr>
<tr>
<td>Zr(( a^3F ))</td>
<td>69</td>
<td>1.00</td>
<td>0.14</td>
<td></td>
</tr>
<tr>
<td>Nb(( a^6D ))</td>
<td>71</td>
<td>1.00</td>
<td>0.04</td>
<td></td>
</tr>
<tr>
<td>Mo*(( a^5S ))</td>
<td>71</td>
<td>1.00</td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

\* \( E_{coll} \) in kJ/mol.

Nh (\( a^6D \)) + 2-Butyne. Both NbC₄H₄ and NbC₃H₂ were observed at \( E_{coll} = 71 \) kJ/mol. The lab angular distributions and TOF spectra were qualitatively similar to those obtained for Zr and are not shown. In fact, for CH₄ elimination, identical CM distributions were used to simulate the Nb and Zr data (Figure 13). Formation of NbC₄H₄ + H₂ was dominant, with \( \phi_{NbC4H4}/\phi_{NbC3H2} \) of 0.04:1.00 (Table 1). To check for reactions of low-lying excited states in the Nb atomic beam, photodepletion studies were performed on the ground \( a^6D \) state using REMPI. Concurrent with a 24 (±4)% integrated depletion of the Nb atomic beam was a 22 (±4) and 23 (±3)% depletion of the NbC₃H₂ and NbC₄H₄ signals, respectively, demonstrating that the reactions primarily involved ground-state Nb atoms.

Figure 11. CM distributions for competing product channels used to generate fits to the data resulting from collisions of Y + 2-butyne at \( E_{coll} = 69 \) kJ/mol shown in Figure 10.
Mo \((a^7S)\) and Mo* \((a^5S)\) + 2-Butyne. For both ground-state Mo \((a^7S)\) and excited-state Mo* \((a^5S)\), only \(H_2\) elimination was observed from reactions with 2-butyne. Simulations for \(H_2\) elimination were qualitatively similar to those obtained for Zr and Nb.

Discussion

Nonreactive Scattering. As already noted, nonreactive scattering has been modeled as involving two processes. In physical terms, species scattered in the forward direction \((\theta \approx 0^\circ)\) resulted from collisions that did not involve complex formation. These collisions are nearly elastic, supporting the notion of minimal interaction between propyne or 2-butyne and the metal atom in these scattering events. The forward-backward symmetric contribution is a result of the decay of long-lived \(\pi\)-association complexes associated with the addition of the \(\text{C} \equiv \text{C}\) triple bond to the metal.\(^{31}\)

The extent of forward scattering relative to the decay of a long-lived \(\pi\)-complex is related to the ability of the ground electronic-state metal atom reactants to form the bound \(\text{M} \equiv \text{propane}\) or \(\text{M} \equiv \text{butyne}\) intermediate. The fraction of complexes decaying back to reactants also depends on the ability of these complexes to go on to reaction, leading to the elimination of a small molecule. The addition of the ethyne \(\text{C} \equiv \text{C}\) bond to second-row metal atoms has been investigated theoretically by Siegbahn.\(^{33}\) For each system, the \(\text{M} \equiv \text{ethyne}\) complex has a different metal electronic configuration than the separated reactants. In Figure 14, the theoretical results obtained by Siegbahn for ground-state complexes have been adapted for reactions of 2-butyne by making the reasonable assumption that the \(\pi\)-complexes involving \(\text{M} + 2\)-butyne are similar to those for \(\text{M} + \text{ethyne}\). The curves shown in Figure 14 representing excited-state longer range metal–alkyne complexes were estimated based on the following considerations: For Y \((\text{doublet, } 4\text{d}^5\text{s}^2)\) and Zr \((\text{triplet, } 4\text{d}^2\text{f}^2)\), promotion of a \(5\text{s}\) electron to a \(4\text{d}\) (or \(5\text{p}\)) orbital is required to form the ground-state \(\text{M} \equiv \text{alkyne}\) complex. For Nb \((\text{sextet, } 4\text{d}^4\text{f}^5\text{h}^1)\) and Mo \((\text{septet, } 4\text{d}^5\text{f}^6\text{i}^1)\), a \(4\text{d}\) electron must undergo a spin flip to form the ground-state complex with an alkyne.\(^{33}\) Theoretical calculations on reactions involving alkenes can be used to provide insight into the nature of the excited-state metal–alkyne complexes. For example, the addition of Y \((\text{a}^2\text{D})\) to \(\text{C}_2\text{H}_4\) can lead to the formation of a long-range \(2\text{B}_2\) complex bound by 69 kJ/mol, or a \(2\text{A}_1\) complex bound by 93 kJ/mol.\(^{34,35}\) For \(\text{Y} + \text{C}_2\text{H}_4\), a long-range \(2\text{B}_2\) complex should be bound by approximately 170 kJ/mol.\(^{35}\) For Zr \((\text{a}^3\text{F}) + \text{C}_2\text{H}_4\), a long-range \(3\text{B}_1\) complex was found to be bound by 40.6 kJ/mol, as compared to 118.0 kJ/mol for the \(\text{ZrC}_2\text{H}_4\) \((3\text{A}_2)\) complex.\(^{36}\) Assuming a similar situation for \(\text{Zr} + \text{C}_2\text{H}_2\) places a long-range \(3\text{B}_1\) complex at about 170 kJ/mol below the reactants, as compared to 238 kcal/mol for \(\text{ZrC}_2\text{H}_4\) \((3\text{A}_2)\).\(^{33}\) Inspection of the electronic configuration, however, suggests that for alkyne, the long-range \(6\text{A}_1\) complex will be destabilized relative to the ground state \(4\text{A}_2\) due to a repulsive interaction between a singly occupied \(\text{d}\) orbital and the out-of-plane alkyne \(\pi\)-bonding orbital. Thus, the long-range \(6\text{A}_1\) alkyne complex is anticipated to lie 84–146 kJ/mol above the ground-state \(4\text{A}_2\) complex. Finally, for Mo, a septet complex with \(\text{C}_2\text{H}_4\) was found to lie 4–8 kJ/mol below reactants.\(^{4}\) For alkyne, a weakly bound septet complex will likely be destabilized by an occupied \(\text{d}\) orbital interacting with the out-of-plane \(\pi\)-system.
The absence of wide-angle scattering for ground-state Mo (αS) atoms indicates that the cross-section for complex formation, which requires intersystem crossing to the quintet surface, is very small. For Y, on the other hand, the cross-section for initial complex formation is very large due to the strongly bound nature of the long-range \(^3\)B\(_2\) complex. While a fraction of these evolve to the \(^3\)A\(_2\) surface and react, a substantial fraction decays back to reactants. The small percentage of symmetric nonreactive scattering for Zr is due to the weaker binding of the long-range \(^1\)B\(_1\) complex and efficient crossing to the more strongly bound reactive \(^3\)A\(_2\) surface where the small C–H bond-activation barrier favors decay to products. For Nb, a larger fraction of complexes decay back to reactants because intersystem crossing to the low-spin \(^3\)A\(_2\) surface is necessary for the reaction to occur. For Mo* (αS), the bound \(^3\)B\(_2\) Mo–propyne complex is readily accessed and can go directly on to form products.

**C–H and C–C Insertion.** Once formed, the M–propyne intermediate may decay back to reactants or undergo insertion. Three possible insertion processes are possible: (1) αC–H insertion into the sp hybridized C–H bond, (2) C–C insertion into the sp–sp\(^3\) hybridized C–C bond, or (3) βC–H insertion into the sp\(^3\) hybridized C–H bonds of the methyl group. Siegbahn calculated the C–H bond insertion barriers for all second-row transition metal atoms with ethyne.\(^{33}\) We present here TS energies that have been modified in accord with more recent calculations.\(^{27}\) Thus, for Y, Zr, Nb, and Mo reactions with propyne, the α-C–H bond insertion barriers are \(-33.0, -31.8, 14.2,\) and \(74.1\) kJ/mol relative to ground-state reactants, respectively.\(^{27}\) While the barriers for β-C–H insertion in these systems have not been calculated, our studies of reactions involving propene\(^{10}\) and four butenes\(^{8}\) have provided compelling evidence that β-C–H insertion barriers into sp hybridized C–H bonds of methyl substituted alkenes are substantially lower than for α-C–H insertion. This explains as to why, for example, in previous room-temperature studies of reactions involving transition metal atoms with alkenes, propene was found to be nearly an order of magnitude more reactive than ethene.\(^{4,31}\) Although the probability of the β-H atoms may be less optimal for alkenes than for alkynes, because M–alkyne binding energies are substantially greater than in M–alkene systems,\(^{3,4}\) for reactions of Y, Zr, and Nb with propyne and 2-butyne, it seems certain that the barriers for β-C–H insertion will lie well below the energy of the separated M + alkene reactants.

To date, there have been no calculations of the potential energy barriers for C–C insertion of metal atoms in alkene and alkyn reactions. However, for reactions of saturated hydrocarbons such as ethane, barrier heights for C–C insertion are larger than C–H bond insertion barriers by 40–80 kJ/mol for neutral metal atom reactions.\(^{3,4}\) Thus, C–C insertion is generally not competitive with C–H insertion, except for systems such as cyclopropane,\(^{8,10}\) where C–C ring strain leads to a substantial decrease in barrier heights; in some cases (e.g., Mo), the calculated C–C insertion barriers actually lie lower than those for C–H insertion.\(^{9,10}\)

In previous studies of reactions of second-row transition metal atoms with propyne,\(^{9,10}\) although C–C activation leading to the production of MCH\(_2\) + CH\(_2\)H\(_2\) was observed, the reaction mechanism did not involve C–C insertion. Instead, the reaction was initiated by π-complex formation, followed by β-C–H insertion and H atom migration, ultimately leading to the formation of a four-membered cyclic metalloccyclobutane complex that decomposes to MCH\(_2\) + CH\(_2\)H\(_2\).\(^{9,10}\) Reactions of propyne and 2-butyne are likely initiated by addition to the C–C triple bond, in these cases forming metalloccyclopropene complexes. Although subsequent β-C–H insertion may occur, the four-membered cyclic intermediate in this case will contain a C=C bond (metallacyclobutene). Because of ring strain, the formation of such intermediates is not as thermodynamically favorable as for metallacyclobutenes. To better assess the role of H atom migration in alkyn reactions, we carried out a limited study of the reaction of Y + allene. Allene (CH\(_2\)–C=CH\(_2\)) is a structural isomer of propyne (CH\(_2\)CCH\(_3\)). The enthalpies of formation of both isomers are nearly identical,\(^{37}\) and both reactions should be initiated by the formation of π-association complexes. At a collision energy of 84 kJ/mol, the reaction dynamics and branching ratios between YC\(_2\)H + CH\(_2\) and YCH\(_2\)H\(_2\) + H are essentially the same for propyne and allene. This strongly suggests that following the formation of the π-association complex, H atom migration is facile.

**H\(_2\) Elimination.** For reactions with propyne, all metals produce MC\(_3\)H\(_2\) + H\(_2\) products. Concerted H\(_2\) elimination from the CH\(_2\) moiety in the initially formed π-association complex is expected to involve a large potential energy barrier and is not expected to play an important role. Thus, H\(_2\) elimination must involve α-C–H insertion, H atom migration, or both. As discussed previously, for propyne, activation of the α-C–H bond (H–CCCH\(_3\)) is favorable for Y, Zr, and Nb. A mechanism connecting the H–M–CCCH\(_3\) intermediate to H\(_2\) elimination involves γ-H migration to form H\(_2\)MCCH\(_3\), which could then undergo H\(_2\) elimination forming MCCH\(_2\) + H\(_2\). However, as already noted, β-C–H bond insertion producing an HCCCH\(_3\)–M–H intermediate is also likely to be favorable. This complex may then undergo either α- or γ-H migration to form M(HCCCH) or MCCCH\(_2\) products, respectively. Despite the fact that the propargylene (HCCCH) isomer has been calculated to be 134.7 kJ/mol more stable than vinylidenecarbene (CCCH\(_2\)),\(^{38}\) we expect the M–CCCH\(_2\) binding energies to be much stronger than those for M–HCCCH\(_3\), resulting in similar energetics for the formation of these two isomers. In summary, it seems likely that both α- and β-C–H bond-activation mechanisms can occur and that both MC\(_3\)H\(_2\) product isomers can be formed.

In the preceding paragraphs, two alternative sequential mechanisms for H\(_2\) elimination were suggested. The strong similarity between the dynamics of Y reactions with propyne and allene provide strong evidence that H atom migration is indeed facile. However, theoretical calculations on a number of simple M–hydrocarbon systems have shown that H\(_2\) elimination frequently proceeds by concerted mechanisms involving multicentered transition states (MCTS),\(^{34,36,40}\) For M + propyne and M + 2-butyne, a larger translational energy release was observed for YC\(_2\)H\(_2\) products than for MC\(_3\)H\(_2\) (M = Zr, Nb, Mo, or Mo*). Similar behavior was observed for other systems\(^{9,8,21,31}\) and provides strong evidence that the barrier for H\(_2\) elimination is large in reactions with yttrium. While no calculations have been performed for M + propyne, it seems likely that there MCTS exist for concerted H\(_2\) elimination that may in fact lie below those for sequential processes.

For reactions with 2-butyne, following formation of the π-complex, the metal atom can activate a β-C–H bond, leading to H–M–CH\(_2\)=CCCH\(_3\). A mechanism for H\(_2\) elimination from this intermediate involves α-H migration, eventually leading to M=CHCCCH\(_3\). Using \(ΔH_f (HCCCH\(_3\)) ≈ 502\) kJ/mol,\(^{24}\) and assuming \(D_h (M=CHCCCH\(_3\)) ≈ D_h (M=CH\(_2\))\), this channel should be \(-63\) kJ/mol endoergic for Y and Mo, and roughly thermoneutral for Zr and Nb. This is not consistent with the \(P(E)\) distributions required to fit the \(E_{coll}\) \(≈ 69\) kJ/mol data, which all contain high-energy tails extending out to 120–160
kJ/mol (for Y, see Figure 11). Instead, formation of a metal diethyne complex, M(C_2H_2)_2, should be quite exoergic (≈100 kJ/mol for Y), consistent with the P(E) distributions in Figure 11 and observed for the other metals.

**CH_3 Elimination.** For reactions with propyne, methyl elimination is only observed for reactions of Y (a^2D) and then only at collision energies of 69 kJ/mol and above. A thermodynamic threshold of 63 ± 8 kJ/mol has been inferred from the data analysis. We have previously observed an analogous channel for reactions of Y (a^2D) + ethyne (elimination of H atoms), where we determined ΔG(ethylene) = 461.1 ± 8.4 kJ/mol. Using this measured bond energy along with known thermodynamic values, assuming that the internal energy of the propyne reactant is zero, the reaction endoergicity for YC_3H + CH_3 can be calculated to be 56.5 ± 8.4 kJ/mol, in good agreement with the threshold observed in this work. Our finding that the observed threshold coincides with the thermodynamic threshold suggests that CH_3 elimination occurs without any significant barrier above the reaction endoergicity.

The absence of CH_3 elimination for reactions of propyne with Zr, Nb, and Mo is probably a consequence of the weak M—CCH bond strengths for these elements as compared to that for Y—CCH. Although no experimental or calculated bond energies are available for these systems, it is probable that these channels are energetically closed even at the highest collision energies studied. For Mo*, the substantial electronic energy makes CH_3 elimination energetically feasible, yet this process was not observed. There are two mechanistic pathways to CH_3 elimination in the Y + propyne and 2-butyne reactions. The first step is the formation of the Y-alkyne complex. Simple C—C bond cleavage can then occur with a loss of CH_3 directly via a barrierless dissociation. Alternatively, insertion into an sp—sp^3 C—C bond forming a R—C≡C—Y—CH_3 intermediate may occur, followed by CH_3 loss. Because both pathways correspond to barrierless elimination of CH_3, our data cannot be used to discriminate between the two possible mechanisms for CH_3 elimination.

For Y + propyne and 2-butyne, as the collision energy was increased, a striking rise in the amount of CH_3 elimination was observed. From previous studies of Y reactions, elimination of H_2 is expected to proceed via a tight TS (i.e., over a potential energy barrier). A similar situation is expected for CH_3 elimination in the case of 2-butyne. Elimination of a radical such as CH_3, however, should proceed via a loose TS (i.e., no potential energy barrier above the product asymptote). In a statistical model, the loose TS for CH_3 elimination leads to an increased number of open states as compared to the tight TS for H_2 and CH_4 elimination. Thus, as the collision energy is increased, the reaction rate increases more rapidly for the elimination of CH_3 than for the elimination of H_2 or CH_4, allowing CH_3 elimination to become more competitive at higher collision energies.

**CH_4 Elimination.** The reaction M + 2-butyne → MC_3H + CH_4 presents a highly exothermic product channel involving C—C bond cleavage. As far as we know, this is the first observation of methane elimination in neutral transition metal—hydrocarbon reactions. For Y, this channel lies approximately 140 kJ/mol below the separated reactants, which is about 40 kJ/mol lower than the products from H_2 elimination.

It is clear that H atom migration must occur for this channel to be operative. As already noted, formal C—C insertion is not necessarily a prerequisite for C—C bond fission. One interesting aspect of the M + propyne and M + 2-butyne reactions is that they allow a comparison of H_2 elimination and CH_4 elimination, respectively, accompanied by the formation of the same product, MC_3H_2. In Figure 15, we present scaled translational energy release distributions, $P(f_T)$, for H_2 elimination (dashed lines) and CH_4 elimination (solid lines) for YC_3H_2 and ZrC_3H_2 formed from the reaction with propyne and 2-butyne, respectively. For both metals, two distributions are included for H_2 elimination at different total energies that bracket the total energy for CH_4 elimination (e.g., $E_{tot} = 138$ and 188 kJ/mol for ZrC_3H_2 + H_2 and $E_{tot} = 176$ kJ/mol for ZrC_3H_2 + CH_4). The distributions for Zr are notably similar, with both H_2 and CH_4 elimination $P(f_T)$ peaking near $f_T = 0.1$. The small $f_T$ values for these distributions are consistent with the absence of a significant potential energy barrier in the exit channel (other than centrifugal barriers). Similar $P(f_T)$ values have been observed experimentally for Zr + CH_2=C=CH_2.

For Y, the $P(E)$ values for the elimination of H_2 and CH_4 show that a greater fraction of the available energy was partitioned into translational energy than for the Zr reactions, suggesting a potential energy barrier in the exit channel for both H_2 and CH_4 elimination in the Y reaction.

As shown in Figure 15, the $P(f_T)$ distribution for CH_4 elimination from Y + 2-butyne peaks further away from the zero of translational energy than those for H_2 elimination from Y + propyne. Thus, a larger amount of translational energy is observed in the system with more product vibrational degrees of freedom. This indicates that for the formation of a given unsaturated metal species involving Y, the exit channel barrier for CH_4 elimination is larger than that for H_2 elimination. This is consistent with expectations for the microscopic reverse process (i.e., that insertion into C—H bonds involves barriers larger than for insertion into H_2). Such behavior may be easily understood by considering the directionality of the participating bonds. In the case of H_2, the spherical s orbitals can participate in multicentered bonding, thereby stabilizing the transition state. For CH_4, on the other hand, the highly directional sp^3 hybridized C—H bond is nearly broken near the top of the barrier.

**C—C versus C—H Activation.** For propyne reactions, competition between C—C and C—H bond activation was only seen for yttrium. The observed product branching ratio, $\phi_{YCH_3}/\phi_{YCH_2}$, was 0.11:1.00 at $E_{coll} = 99$ kJ/mol and decreased at lower collision energies. For M + 2-butyne, the competition between C—H and C—C bond activation can be understood by comparing the product branching ratios (Table 1). For Y, C—C bond activation (as both CH_3 and CH_4 elimination) constituted as much as 27% of the total products at $E_{coll} = 116$ kJ/mol. Elimination of CH_4 was also significant for Zr, albeit less so than for Y. The NbC_3H_2 + CH_4 channel was observed at $E_{coll}$...
observed, even at high collision energies. For Y-...the series propane, propene, and propyne, and for alkanes, alkenes, and alkynes indicates that C-...elimination of H2. This potential energy barrier in the exit channel for the most...unique among the metal atoms studied. One major factor that contributes to this rich chemistry of Y is the presence of a potential energy barrier in the exit channel for the most thermodynamically favored channel, the elimination of H2. This allows less thermodynamically favorable products to be competitive. For butyne reactions, this factor likely also plays an important role in facilitating competition with dynamically less favorable channels.

Conclusion

The collisions of early second-row transition metal atoms with propyne and 2-butyne were studied using crossed molecular beams. The reaction M + propyne → MC3H2 + H2...and NbC3H2 were similar in shape for these two reactions, consistent with both H2 and CH4 elimination involving a loose transition state. For Y, considerably more translational energy was released than for Zr or Nb, suggesting a potential energy barrier in the exit channels.

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(42) CM distributions for Zr and Nb were very similar for these reactions; thus, similar considerations as those discussed for Zr would be expected to hold for Nb as well.